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### **LASER METROLOGY OF ANISOTROPIC POLYMER LAYERS STRUCTURE OF MATERIALS FOR PACKAGING INDUSTRY**

**We have proposed the use of a laser-polarimetric method, as more informative in terms of displaying optically inhomogeneous structures by using spatially coherent filtering.**

**Keywords: applied programming; polymers; publishing; digital and graphic information; packaging materials; materials science.**

#### **Introduction**

The main idea of using filtering [1–5] is that the spatial-frequency nature of the Fourier images of the polymer layer laser image of the packaging material is different in its large-scale and small-scale components.

Thus, spatial-frequency filtering can isolate both low-frequency and high-frequency components, which are then transformed into the corresponding filtered images of multiscale anisotropic structures using the inverse Fourier transform [6, 7].

The article provides a substantiation of the principles of consistent spatial-frequency polarization-correlation selection for investigation the linear and circular phase anisotropy of polymeric materials.

#### **Method**

The analysis of the conversion of laser radiation by polycrystalline grids of various polymeric spheres is based on the following model representations [8–10]:

— the polymer is considered as a two-component isotropic-anisotropic (birefringent) structure;



— the optical-anisotropic component of the polymer is formed by linearly birefringent whiskers, which introduce different spatial-frequency phase modulation between the orthogonal components of the laser beam amplitude.

If a low-pass filter  $R$  or a high-pass filter  $R^{-1}$  is placed in the central part of the Fourier plane



$$R(\Delta\eta, \Delta\mu) = \begin{cases} 1 \rightarrow \eta^* \in \Delta\eta; \mu^* \in \Delta\mu; \\ 0 \rightarrow \eta^* \notin \Delta\eta; \mu^* \notin \Delta\mu, \end{cases} \quad (1)$$



$$R^{-1}(\Delta\eta, \Delta\mu) = \begin{cases} 1 \rightarrow \eta^* \notin \Delta\eta; \mu^* \notin \Delta\mu; \\ 0 \rightarrow \eta^* \in \Delta\eta; \mu^* \in \Delta\mu, \end{cases} \quad (2)$$

after it is possible to distinguish the next spatial-frequency structures of the Fourier spectra of the azimuth and ellipticity distributions of the boundary field, that formed either mainly by the effects of linear  $\hat{U}(\kappa, \gamma, \eta, \mu)$  or circular  $\hat{U}(\psi, \eta, \mu)$  anisotropy

$$\begin{cases} \hat{U}(\kappa, \gamma, \eta, \mu) = \\ = R(\Delta\eta, \Delta\mu)U(\eta, \mu); \\ \hat{U}(\psi, \eta, \mu) = \\ = R^{-1}(\Delta\eta, \Delta\mu)U(\eta, \mu). \end{cases} \quad (3)$$

The corresponding distributions of complex amplitudes in the image plane of the polymer layer can be restored using the inverse FFT

$$\begin{bmatrix} \hat{E}_x(\kappa, \gamma, x, y) \\ \hat{E}_x(\psi, x, y) \end{bmatrix} \text{ and } \begin{bmatrix} \hat{E}_y(\kappa, \gamma, x, y) \\ \hat{E}_y(\psi, x, y) \end{bmatrix}.$$

$$\begin{cases} \hat{E}_x(\kappa, \gamma, x, y) = \\ = \frac{1}{i\lambda f} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} R(\Delta\eta, \Delta\mu) \cdot \\ \cdot \hat{U}_x(\eta, \mu) \exp \cdot \\ \cdot [i2\pi(x\eta + y\mu)] d\eta d\mu \\ \hat{E}_x(\psi, x, y) = \\ = \frac{1}{i\lambda f} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} R^{-1}(\Delta\eta, \Delta\mu) \cdot \\ \cdot \hat{U}_x(\eta, \mu) \exp \cdot \\ \cdot [i2\pi(x\eta + y\mu)] d\eta d\mu \end{cases} \quad (4)$$

$$\begin{cases} \hat{E}_y(\kappa, \gamma, x, y) = \\ = \frac{1}{i\lambda f} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} R(\Delta\eta, \Delta\mu) \cdot \\ \cdot \hat{U}_y(\eta, \mu) \exp \cdot \\ \cdot [i2\pi(x\eta + y\mu)] d\eta d\mu \\ \hat{E}_y(\psi, x, y) = \\ = \frac{1}{i\lambda f} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} R^{-1}(\Delta\eta, \Delta\mu) \cdot \\ \cdot \hat{U}_y(\eta, \mu) \exp \cdot \\ \cdot [i2\pi(x\eta + y\mu)] d\eta d\mu \end{cases} \quad (5)$$

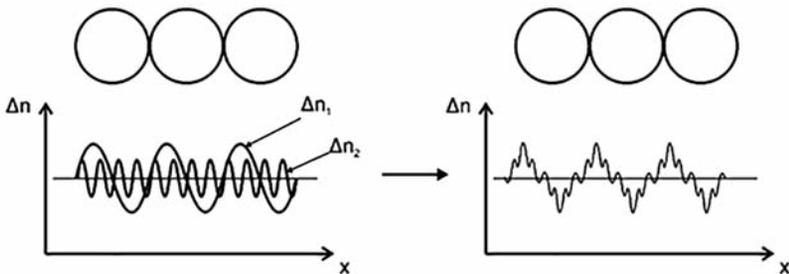


Fig. 1. To the analysis of model representations (object A)

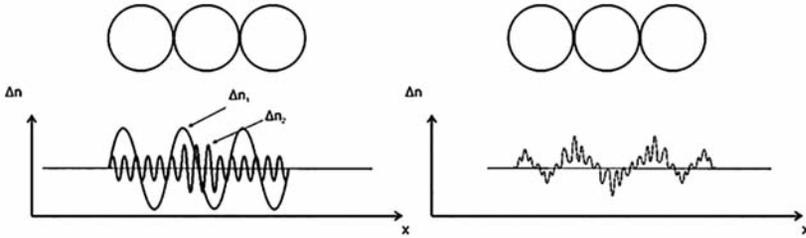


Fig. 2. To the analysis of model representations (object — type B)

Taking into deviation (4) and (5), polarization azimuth maps are determined for the next spatial-frequency filtering

$$\begin{cases} \hat{\alpha}(\kappa, \gamma) = \\ = 0,5 \arctg \left( \frac{\hat{E}_x \hat{E}_y^* - \hat{E}_y \hat{E}_x^*}{\hat{E}_x \hat{E}_x^* - \hat{E}_y \hat{E}_y^*} \right) \\ \hat{\alpha}(\psi) = \\ = 0,5 \arctg \left( \frac{\hat{E}_x \hat{E}_y^* - \hat{E}_y \hat{E}_x^*}{\hat{E}_x \hat{E}_x^* - \hat{E}_y \hat{E}_y^*} \right) \end{cases} \quad (6)$$

**Results**

We use two types of model polymer objects:

1. Object 1 (A) — it is a network of ordered circular cylinders with two-frequency harmoniously components (low-frequency component

$$\Delta n_1 = \Delta n_0 \sin\left(\frac{2\pi}{4X_0} x\right)$$

and high-frequency component  $\Delta n_2 = \Delta n_0^* \sin\left(\frac{2\pi}{X_0} x\right)$

linear anisotropy (fig. 1);

2. Object 2 (B) — it's a similar network, where the central five cylinders are characterized by an increased level (up to 3 times) of the linear birefringence  $\Delta n_0^*$  (fig. 2).

On fig. 3 shows the coordinate distributions of the maps of azimuth of polarization of the Fourier spectrum of the type A network (a) and the spatially-frequency filtered distribution of the maps of azimuth (b).

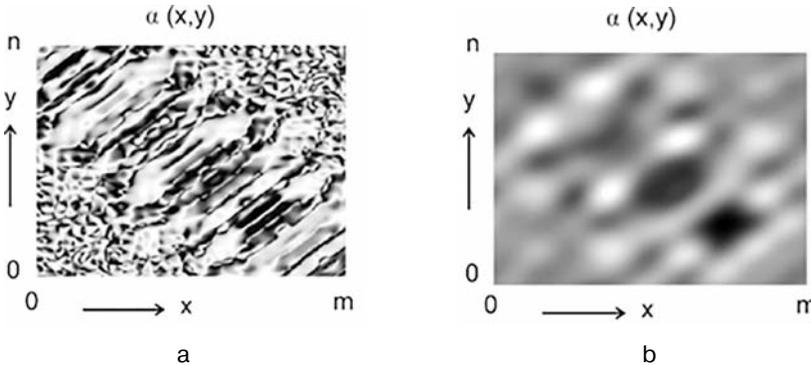


Fig. 3. Coordinate distributions (a) and large-scale (b) component of the azimuth of polarization (type A)



On fig. 4 are shown the spatial-frequency filtered distributions of the azimuths maps of the large-scale component of the polycrystalline polymer network with linear birefringence (a) and their statistical characteristics (b).

Comparative characteristics of polarization azimuth maps are illustrated in table 1.

Table 1  
Statistical moments of a large-scale polarization map of the azimuth

Parameters	LS (large-scale)	
	A-type	B-type
$M_1$	0.29	0.33
$M_1$	0.23	0.27
$M_1$	0.83	1.69
$M_1$	1.13	3.47

**Discussion**

On fig. 5 shows the coordinate distributions of the azimuth maps

of a model object (type A) before and after using spatial-frequency filtering with a high-frequency  $\Delta R^{-1} = = 60$  pixels (size of filter).

The data obtained demonstrate the possibility of using a high-frequency filter, which makes it possible to effectively extract the small-scale component of the coordinate distribution of azimuths by using the inverse Fourier transform.

The results of statistical analysis are presented in fig. 6 and in table 2.

Table 2  
Statistical moments of a small-scale map of the azimuth

Parameters	HF (small-scale)	
	A-type	B-type
$M_1$	0.074	0.087
$M_1$	0.12	0.14
$M_1$	0.56	0.81
$M_1$	2.06	0.58

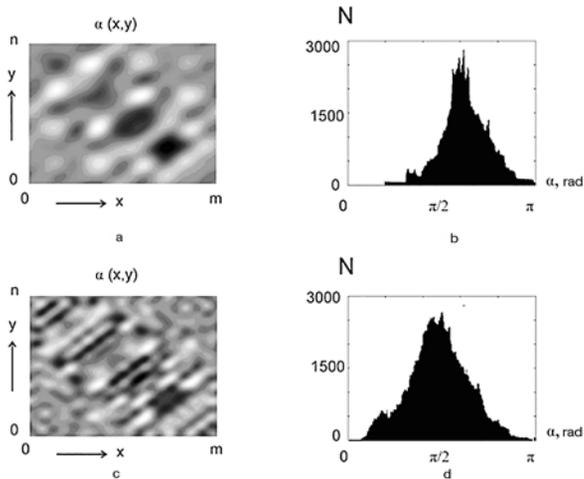


Fig. 4. Large-scale map (a, c) and histogram (b, d) of the coordinate distribution of the maps of azimuth (A-type of polycrystalline network)

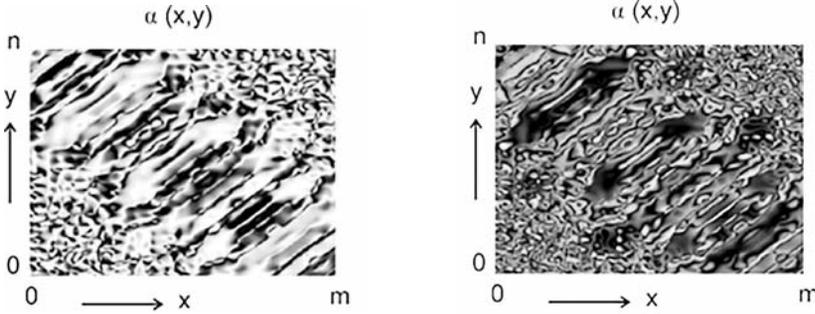


Fig. 5. Coordinate distribution ( $m \cdot n$ ) and their small-scale component of the azimuth maps of both types (A and B networks)

Comparison of the parameters that characterize the coordinate distributions of the polarization azimuth showed the following main changes in the high-frequency component of linear birefringence.

Firstly, this is a significant expansion of the values, as well as an increase in the histogram scatter of the small-scale component of the polarization azimuth  $\alpha(x, y)$  maps, which is formed by the B-type model grid.

Secondly, an increase in the depth of modulation of the high-frequen-

cy component of linear birefringence manifests itself in the formation of an inhomogeneous distribution  $\alpha(x, y)$ .

**Conclusions**

Comparison of the obtained results of spatial-frequency filtering of the distributions of polarization azimuths of model objects revealed the following parameters, which are diagnostically sensitive to changes in birefringence:

— statistical moments of the 3rd and 4th order of orders. The differ-

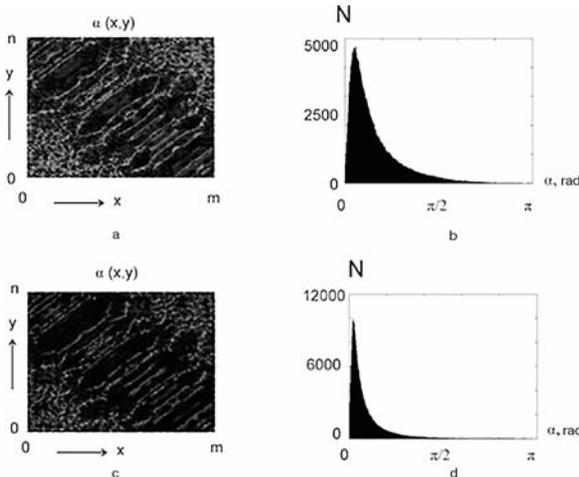


Fig. 6. Small-scale maps (a, c) and histograms (b, d) of the azimuth distributions of a polycrystalline networks (A-type and B-type)



ences between the values of this data of both types range from 2 ( $M_3$ ) to 3.5 ( $M_4$ ) times;

— statistical moments of the 2nd–4th orders of the azimuth distribu-

tions and small-scale component of the polarization map of the image of a grid of optically birefringence cylinders. The differences from 1.5 ( $M_2$ ,  $M_3$ ) to 4 ( $M_4$ ) times.

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**В роботі запропоновано використання лазерно-поляриметричного методу, який є більш інформативний в сенсі відображення оптично неоднорідних структур за рахунок використання просторово-когерентної фільтрації.**

**Ключові слова: прикладне програмування; полімери; пакувальна продукція; цифрова та графічна інформація; матеріалознавство.**

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